## Generation of increasing light signals during the parametrical process on photo-integrated anisotropy

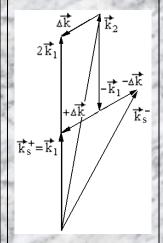
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The observation results of the generation of increasing light signals during the parametrical process on photo-integrated anisotropy are presented. The sufficiently big growth of the intensity of small signals of the basic frequency infra-red radiation with time (photo-stimulated amplification of light) has been detected in experiments with the influence on amorphous optical material by the weak probe radiation of the basic frequency in presence simultaneously the strong radiation of the second harmonic. Various modes of the amplification of light signals were investigated on preliminary created photo-integrated space-periodical anisotropy with different initial amplitudes. The results of the detailed studies of the generation of increasing light signals on photo-integrated anisotropy have been presented with demonstration of the investigated properties of the observed parametrical process in dependence on intensities of the pumping radiation, on polarization of the light signals, on the phase shifts and on the spatial distributions. Mechanisms of the occurrence of the parametrical generation process on photo-integrated anisotropy with the existence of the non-steady regimes of signal amplification are also discussed. The presented investigations may be interesting for the researchers in a spheres of photonics and micro-optoelectronics. The work was carried out as the part of tasks of the Russian State Project FWGW-2021-0012.

For nonlinear process of parametrical down frequency conversion (or parametric amplification of low light signals) the photo-integrated by volumetric all-optical poling susceptibility tensor of  $\chi_{ini}^{(2)}$  grating:

 $\chi_{ijk}^{(2)}(\Omega; 2\Omega, -\Omega) = (\chi_1^{(3)} \boldsymbol{E}_k \delta_{ij} + \chi_2^{(3)} \boldsymbol{E}_j \delta_{ik} + \chi_3^{(3)} \boldsymbol{E}_i \delta_{kj}),$ 

 $\chi_1^{(3)} = 2\pi\chi_{iiji}^{(3)}(\Omega; 2\Omega, -\Omega), \quad \chi_2^{(3)} = 2\pi\chi_{iji}^{(3)}(\Omega; 2\Omega, -\Omega), \quad \chi_3^{(3)} = 2\pi\chi_{iji}^{(3)}(\Omega; 2\Omega, -\Omega).$ So, maximal nonlinear conversion of light observed for radiation polarized along the direction of the induced electric field *E*. But there are the other non-zero components.



conversion of light observed for radiation polarized along the electric field E. But there are the other non-zero components. Fig 1. Phase synchronism on photo-induced grating  $\chi_{int}^{(2)}$  for nonlinear process of degenerate parametrical down frequency conversion with amplification of low light signals of basic frequency. The low signal wave with the wave vector  $\mathbf{k}_s = \mathbf{k}_1$  and the big power pumping wave with the vector  $\mathbf{k}_p = \mathbf{k}_2$  interact in medium with photo-integrated  $\chi_{int}^{(2)}$  grating with vector  $\Delta \mathbf{k} = 2\mathbf{k}_1 - \mathbf{k}_2$  by condition of phase synchronization  $\mathbf{k}_s = \mathbf{k}_2 - \mathbf{k}_1 \pm \Delta \mathbf{k}$ . As the  $\mathbf{k}_s^+$  coincide with incident basic light beam  $\mathbf{k}_1$ , the spatial phase synchronism for inducing  $\chi_{int}^{(2)}$  grating and signal wave with basic frequency is carried out automatically and there is a case of the nonlinear degenerate parametrical amplification of low signals.

