

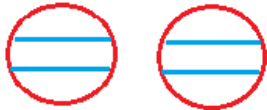
Dynamical symmetry of two qubits in external fields with dipole – dipole interaction

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Abstract

Applications of dynamical groups to the temporal dynamics of multilevel quantum systems in external fields are considered. Using the representation of coherent states (CS), it is shown that the time evolution of the state vector is reduced to the "classical" dynamics of complex parameters of the CS, taking values in the space of cosets of the dynamical group of the Hamiltonian. A particular case of the dynamics of two dipole-dipole (dd) interacting qubits used in the schemes of modern quantum informatics is studied in framework of $SU(4)$ group.

Model and approach



ω_1

ω_2

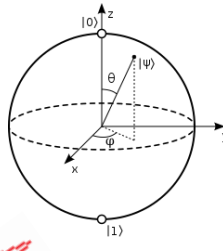
Two dd-interacted qubits

$$\hat{H}(t) = \hat{H}_0 + \hat{V}_I(t), \quad \hat{H}_0 = \hbar\omega_1 \hat{J}_3^{(1)} + \hbar\omega_2 \hat{J}_3^{(2)},$$

$$V_I(t) = -\frac{\hbar g^{(12)}}{2} \left(J_+^{(1)} J_-^{(2)} + J_-^{(1)} J_+^{(2)} \right) -$$

$$-\frac{\hbar}{2} \left\{ \left(\Omega_1 J_+^{(1)} + \Omega_2 J_+^{(2)} \right) \exp(i\Delta\omega t) + h.c. \right\}.$$

$$G = SU(2) \times SU(2) \subset SU(4),$$



$$|z_1, z_2\rangle = |z_1\rangle \otimes |z_2\rangle, \quad |z_\alpha\rangle = (1 + z_\alpha \bar{z}_\alpha)^{-1/2} \exp(z_\alpha \hat{J}_+^{(\alpha)}) |1/2, \downarrow\rangle, \quad \hat{J}_+ = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}.$$

$|z_1\rangle \otimes |z_2\rangle \Rightarrow |\xi\rangle \equiv |\xi_1, \xi_2, \xi_3\rangle - SU(4)$ coherent state;

$$i\hbar \sum_{\beta=1}^3 \frac{\partial^2 \ln K(\xi, \bar{\xi})}{\partial \bar{\xi}^\alpha \partial \xi^\beta} \dot{\xi}^\beta = \frac{\partial \mathcal{H}(\xi, \bar{\xi}; t)}{\partial \bar{\xi}^\alpha} \Leftarrow \text{exact equations!}$$

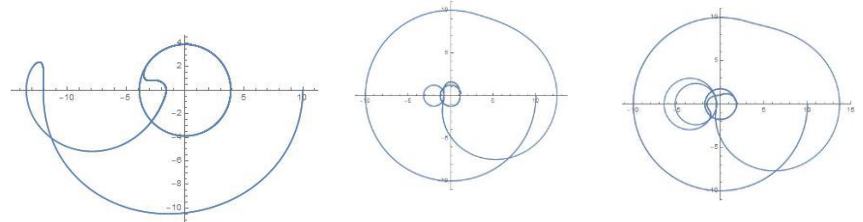
$\alpha, \beta = 1, 2, 3.$

$$\mathcal{H}(\xi, \bar{\xi}; t) = \langle \xi | \hat{H}(t) | \xi \rangle, \quad K(\xi, \bar{\xi}) = \left(1 + \xi_1 \bar{\xi}_1 + \xi_2 \bar{\xi}_2 + \xi_3 \bar{\xi}_3 \right)^{-1}.$$

$$|z_1\rangle \otimes |z_2\rangle = (1 + z_1 \bar{z}_1)^{-1/2} \cdot (1 + z_2 \bar{z}_2)^{-1/2} \begin{pmatrix} z_1 \cdot z_2 \\ z_2 \\ z_1 \\ 1 \end{pmatrix}, \text{ (no qubits entanglement!)}$$

$$|\xi_1, \xi_2, \xi_3\rangle = (1 + \xi_1 \bar{\xi}_1 + \xi_2 \bar{\xi}_2 + \xi_3 \bar{\xi}_3)^{-1/2} \begin{pmatrix} \xi_3 \\ \xi_2 \\ \xi_1 \\ 1 \end{pmatrix}, \quad (\xi_3 \neq \xi_1 \cdot \xi_2 - \text{entanglement!})$$

Numerical calculations



Trajectories in complex planes ξ_1, ξ_2, ξ_3 , respectively.

$$\xi_3(0) = \xi_1(0) \xi_2(0)$$

References

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- 4) Horodecki, R., et al. Quantum entanglement. Rev.Mod.Phys.V. 81. P. 865–942. (2009).

